

Degree-2, Physics (Hons), Paper-B IIIrd
Lecture-5

Maxwell's Equations:→

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0} \quad \text{--- (i) --- Gauss' Law in electrostatics}$$

$$\vec{\nabla} \cdot \vec{B} = 0 \quad \text{--- (ii) --- Gauss' Law for magnetism}$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad \text{--- (iii) --- Faraday's Law}$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t} \quad \text{--- (iv) --- Ampere's Law}$$

where,

$\vec{\nabla}$ → Del Operator

\vec{E} → Electric field intensity.

ρ → charge density

ϵ_0 → Permittivity in free space

\vec{B} → magnetic field

\vec{J} → current density

μ_0 → Permeability in free space.

Integral form of Maxwell's Equations:→

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

Taking volume integral on both side

$$\int \vec{\nabla} \cdot \vec{E} \, dV = \int \frac{\rho}{\epsilon_0} \, dV$$

using Gauss' Divergence theorem

$$\int \vec{\nabla} \cdot \vec{E} \, dV = \int \vec{E} \cdot d\vec{S}$$

$$\boxed{\int \vec{E} \cdot d\vec{S} = \frac{q}{\epsilon_0}}$$

$$(\because q = \int \rho \, dV)$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\int \vec{\nabla} \cdot \vec{B} \, dv = 0$$

$$\boxed{\int \vec{B} \cdot d\vec{S} = 0}$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

Taking surface integral on both side.

$$\int (\vec{\nabla} \times \vec{E}) \cdot d\vec{S} = \int -\frac{\partial \vec{B}}{\partial t} \cdot d\vec{S}$$

using Stokes's theorem,

$$\int (\vec{\nabla} \times \vec{E}) \cdot d\vec{S} = \int \vec{E} \cdot d\vec{l}$$

$$\int \vec{E} \cdot d\vec{l} = -\frac{\partial}{\partial t} \int \vec{B} \cdot d\vec{S}$$

$$\int \vec{E} \cdot d\vec{l} = -\frac{\partial \Phi_B}{\partial t}$$

$$(\Phi_B = \int \vec{B} \cdot d\vec{S})$$

$$\boxed{\int \vec{E} \cdot d\vec{l} = -\frac{\partial \Phi_B}{\partial t}}$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

$$\int (\vec{\nabla} \times \vec{B}) \cdot d\vec{S} = \mu_0 \int \vec{J} \cdot d\vec{S} + \mu_0 \epsilon_0 \frac{\partial}{\partial t} \int \vec{E} \cdot d\vec{S}$$

$$\boxed{\int \vec{B} \cdot d\vec{l} = \mu_0 I + \mu_0 \epsilon_0 \frac{\partial \Phi_E}{\partial t}}$$

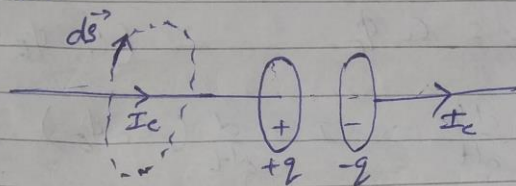
where $I = \int \vec{J} \cdot d\vec{S}$

$$\Phi_E = \int \vec{E} \cdot d\vec{S}$$

Displacement Current \rightarrow

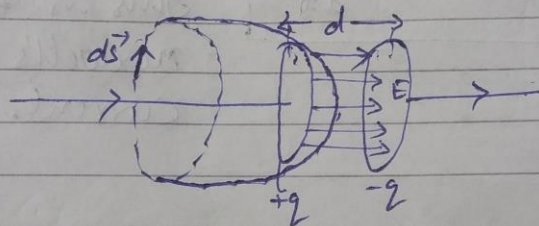
Displacement current is a quantity appearing in Maxwell's equations that is defined in terms of rate of change of electric displacement field.

Consider a charge capacitor. Conducting wires bring current I_c on to one plate and away from the other as the charge on the plates increases.



For above path we can apply Ampere's Law

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I_c$$



The instantaneous charge on the capacitor plates is $q = C \Delta V$, where ΔV is the instantaneous potential difference across the plates.

As capacitor charge the electric field between the plates is changing.

$$C = \epsilon_0 \frac{A}{d}$$

where A is the area of the plate and d is the separation between the plates

and C is the Capacitance of the Capacitor.

$$C = \epsilon_0 \frac{A}{d}$$

$$\Delta V = Ed$$

$$\therefore Q = C \Delta V$$

$$Q = \epsilon_0 \frac{A}{d} \times Ed$$

$$Q = \epsilon_0 AE$$

$$Q = \epsilon_0 \phi_E \quad (\because \phi_E = EA)$$

Now

$$I_c = \frac{dQ}{dt} \Rightarrow I = \frac{d(\epsilon_0 \phi_E)}{dt}$$

$$I_c = \epsilon_0 \frac{d\phi_E}{dt}$$

$$\boxed{I_D = \epsilon_0 \frac{d\phi_E}{dt}}$$

This is known as
Maxwell's displacement
Current.